

**Preliminary Exam: Electromagnetism, Thursday August 24, 2017. 9:00-12:00**

Answer a total of any **THREE** out of the four questions. Use the blue solution books and put the solution to each problem in a separate blue book and put the number of the problem on the front of each blue book. Be sure to put your name on each blue book that you submit. If you submit solutions to more than three problems, only the first three problems as listed on the exam will be graded.

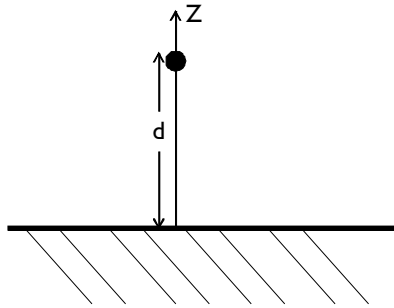
1. A static charge distribution produces a radial electric field of the form:  $E(r) = A \exp(-br)/r$ . The direction of  $\mathbf{E}$  is along the vector  $\mathbf{r}$ , and  $A$  and  $b$  are constants.

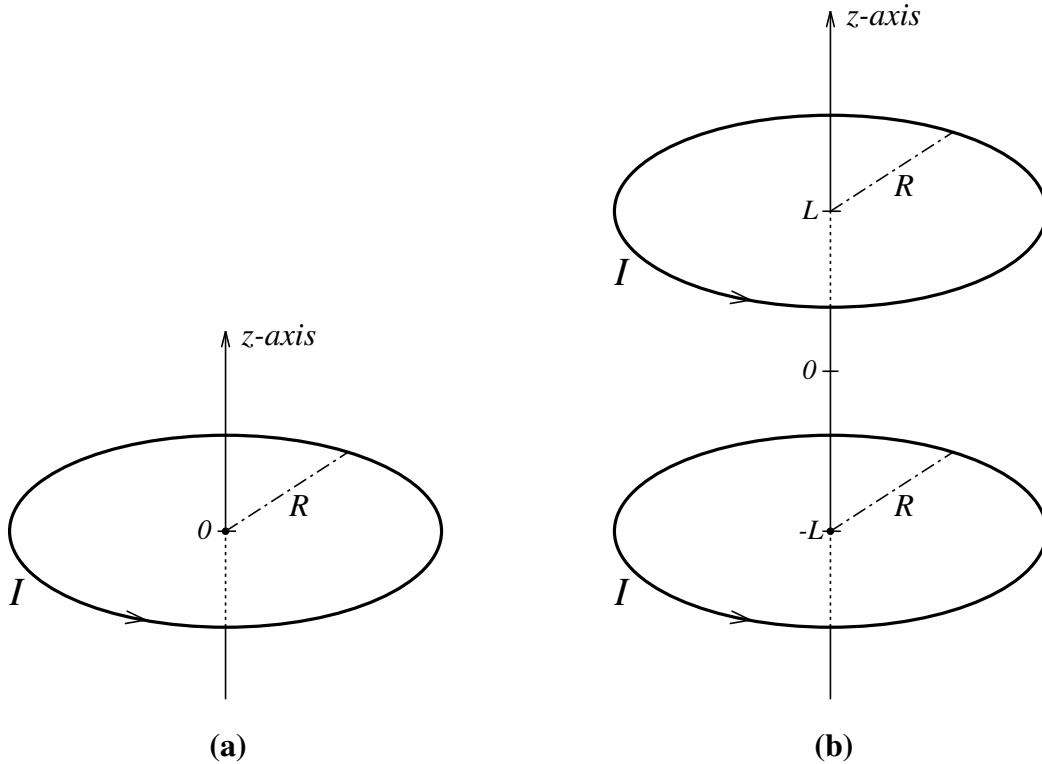
(a) Determine the dependence of the charge density  $\rho(r)$  on  $r$  and draw a sketch of it.

(b) What is the total charge  $Q$ ?

2. (a) Find the work required to bring a charge  $q$  from infinity to a distance  $d$  away from a semi-infinite dielectric medium (see figure) whose dielectric constant is  $K$ .

(b) Calculate the work required if instead the plane is perfectly conducting.





3. (a) Consider a circular loop with radius  $R$  carrying a current  $I$  flowing counter clockwise in the  $x$ - $y$  plane around the origin, see Fig. (a). Determine the magnetic induction along the  $z$ -axis as a function of  $z$ , i.e.  $B_z(\mathbf{x})|_{\mathbf{x}=(0,0,z)} \equiv B(z)$ .
- (b) Consider now two current loops, as shown in Fig. (b), where the distance between the centers of the two loops is  $2L$ . Determine the magnetic induction along the  $z$ -axis.
- (c) Expand the result from part (b) to appropriate order in  $z$  to determine for which  $L$  (expressed in terms of  $R$ ) the magnetic induction along the  $z$ -axis is  $B(z) = B_0 + \mathcal{O}(z^4)$  for  $z \ll L$ . Determine  $B_0$  as function of  $I$  and  $R$ .
4. (a) Consider electromagnetic potentials  $\phi(\mathbf{x}, t)$  and  $\mathbf{A}(\mathbf{x}, t)$  such that the homogeneous Maxwell equations are automatically satisfied. Show that the inhomogeneous Maxwell equations in the presence of a charge density  $\rho(\mathbf{x}, t)$  and a current density  $\mathbf{j}(\mathbf{x}, t)$  decouple as

$$\nabla^2 \phi(\mathbf{x}, t) - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \phi(\mathbf{x}, t) = -\frac{1}{\epsilon_0} \rho(\mathbf{x}, t), \quad \nabla^2 \mathbf{A}(\mathbf{x}, t) - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \mathbf{A}(\mathbf{x}, t) = -\mu_0 \mathbf{j}(\mathbf{x}, t)$$

if one imposes the Lorenz gauge condition  $\nabla \cdot \mathbf{A}(\mathbf{x}, t) + \frac{1}{c^2} \frac{\partial}{\partial t} \phi(\mathbf{x}, t) = 0$  on the potentials. (There are no medium effects in this problem, i.e.  $\epsilon = \epsilon_0$ ,  $\mu = \mu_0$ ,  $\epsilon_0 \mu_0 c^2 = 1$ .)

- (b) Set  $\mathbf{A}(\mathbf{x}, t) = -\frac{z^2}{2} E_0 \frac{\omega}{c^2} \sin(\omega t) \mathbf{e}_z$  in the Lorenz gauge. Determine  $\phi(\mathbf{x}, t)$ ,  $\rho(\mathbf{x}, t)$ ,  $\mathbf{j}(\mathbf{x}, t)$  and check explicitly whether electric charge is conserved. Determine also  $\mathbf{E}(\mathbf{x}, t)$ ,  $\mathbf{B}(\mathbf{x}, t)$  and investigate whether energy is conserved or whether this is an open system.

# Vector Formulas

$$\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = \mathbf{b} \cdot (\mathbf{c} \times \mathbf{a}) = \mathbf{c} \cdot (\mathbf{a} \times \mathbf{b})$$

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \cdot \mathbf{c})\mathbf{b} - (\mathbf{a} \cdot \mathbf{b})\mathbf{c}$$

$$(\mathbf{a} \times \mathbf{b}) \cdot (\mathbf{c} \times \mathbf{d}) = (\mathbf{a} \cdot \mathbf{c})(\mathbf{b} \cdot \mathbf{d}) - (\mathbf{a} \cdot \mathbf{d})(\mathbf{b} \cdot \mathbf{c})$$

$$\nabla \times \nabla \psi = 0$$

$$\nabla \cdot (\nabla \times \mathbf{a}) = 0$$

$$\nabla \times (\nabla \times \mathbf{a}) = \nabla(\nabla \cdot \mathbf{a}) - \nabla^2 \mathbf{a}$$

$$\nabla \cdot (\psi \mathbf{a}) = \mathbf{a} \cdot \nabla \psi + \psi \nabla \cdot \mathbf{a}$$

$$\nabla \times (\psi \mathbf{a}) = \nabla \psi \times \mathbf{a} + \psi \nabla \times \mathbf{a}$$

$$\nabla(\mathbf{a} \cdot \mathbf{b}) = (\mathbf{a} \cdot \nabla)\mathbf{b} + (\mathbf{b} \cdot \nabla)\mathbf{a} + \mathbf{a} \times (\nabla \times \mathbf{b}) + \mathbf{b} \times (\nabla \times \mathbf{a})$$

$$\nabla \cdot (\mathbf{a} \times \mathbf{b}) = \mathbf{b} \cdot (\nabla \times \mathbf{a}) - \mathbf{a} \cdot (\nabla \times \mathbf{b})$$

$$\nabla \times (\mathbf{a} \times \mathbf{b}) = \mathbf{a}(\nabla \cdot \mathbf{b}) - \mathbf{b}(\nabla \cdot \mathbf{a}) + (\mathbf{b} \cdot \nabla)\mathbf{a} - (\mathbf{a} \cdot \nabla)\mathbf{b}$$

If  $\mathbf{x}$  is the coordinate of a point with respect to some origin, with magnitude  $r = |\mathbf{x}|$ ,  $\mathbf{n} = \mathbf{x}/r$  is a unit radial vector, and  $f(r)$  is a well-behaved function of  $r$ , then

$$\nabla \cdot \mathbf{x} = 3$$

$$\nabla \times \mathbf{x} = 0$$

$$\nabla \cdot [\mathbf{n}f(r)] = \frac{2}{r}f + \frac{\partial f}{\partial r} \quad \nabla \times [\mathbf{n}f(r)] = 0$$

$$(\mathbf{a} \cdot \nabla)\mathbf{n}f(r) = \frac{f(r)}{r} [\mathbf{a} - \mathbf{n}(\mathbf{a} \cdot \mathbf{n})] + \mathbf{n}(\mathbf{a} \cdot \mathbf{n}) \frac{\partial f}{\partial r}$$

$$\nabla(\mathbf{x} \cdot \mathbf{a}) = \mathbf{a} + \mathbf{x}(\nabla \cdot \mathbf{a}) + i(\mathbf{L} \times \mathbf{a})$$

where  $\mathbf{L} = \frac{1}{i}(\mathbf{x} \times \nabla)$  is the angular-momentum operator.

# Theorems from Vector Calculus

In the following  $\phi$ ,  $\psi$ , and  $\mathbf{A}$  are well-behaved scalar or vector functions,  $V$  is a three-dimensional volume with volume element  $d^3x$ ,  $S$  is a closed two-dimensional surface bounding  $V$ , with area element  $da$  and unit outward normal  $\mathbf{n}$  at  $da$ .

$$\int_V \nabla \cdot \mathbf{A} d^3x = \int_S \mathbf{A} \cdot \mathbf{n} da \quad (\text{Divergence theorem})$$

$$\int_V \nabla \psi d^3x = \int_S \psi \mathbf{n} da$$

$$\int_V \nabla \times \mathbf{A} d^3x = \int_S \mathbf{n} \times \mathbf{A} da$$

$$\int_V (\phi \nabla^2 \psi + \nabla \phi \cdot \nabla \psi) d^3x = \int_S \phi \mathbf{n} \cdot \nabla \psi da \quad (\text{Green's first identity})$$

$$\int_V (\phi \nabla^2 \psi - \psi \nabla^2 \phi) d^3x = \int_S (\phi \nabla \psi - \psi \nabla \phi) \cdot \mathbf{n} da \quad (\text{Green's theorem})$$

In the following  $S$  is an open surface and  $C$  is the contour bounding it, with line element  $d\mathbf{l}$ . The normal  $\mathbf{n}$  to  $S$  is defined by the right-hand-screw rule in relation to the sense of the line integral around  $C$ .

$$\int_S (\nabla \times \mathbf{A}) \cdot \mathbf{n} da = \oint_C \mathbf{A} \cdot d\mathbf{l} \quad (\text{Stokes's theorem})$$

$$\int_S \mathbf{n} \times \nabla \psi da = \oint_C \psi d\mathbf{l}$$

# Explicit Forms of Vector Operations

Let  $\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3$  be orthogonal unit vectors associated with the coordinate directions specified in the headings on the left, and  $A_1, A_2, A_3$  be the corresponding components of  $\mathbf{A}$ . Then

Cartesian  
( $x_1, x_2, x_3 = x, y, z$ )

$$\begin{aligned}\nabla\psi &= \mathbf{e}_1 \frac{\partial\psi}{\partial x_1} + \mathbf{e}_2 \frac{\partial\psi}{\partial x_2} + \mathbf{e}_3 \frac{\partial\psi}{\partial x_3} \\ \nabla \cdot \mathbf{A} &= \frac{\partial A_1}{\partial x_1} + \frac{\partial A_2}{\partial x_2} + \frac{\partial A_3}{\partial x_3} \\ \nabla \times \mathbf{A} &= \mathbf{e}_1 \left( \frac{\partial A_3}{\partial x_2} - \frac{\partial A_2}{\partial x_3} \right) + \mathbf{e}_2 \left( \frac{\partial A_1}{\partial x_3} - \frac{\partial A_3}{\partial x_1} \right) + \mathbf{e}_3 \left( \frac{\partial A_2}{\partial x_1} - \frac{\partial A_1}{\partial x_2} \right) \\ \nabla^2\psi &= \frac{\partial^2\psi}{\partial x_1^2} + \frac{\partial^2\psi}{\partial x_2^2} + \frac{\partial^2\psi}{\partial x_3^2}\end{aligned}$$

Cylindrical  
( $\rho, \phi, z$ )

$$\begin{aligned}\nabla\psi &= \mathbf{e}_1 \frac{\partial\psi}{\partial\rho} + \mathbf{e}_2 \frac{1}{\rho} \frac{\partial\psi}{\partial\phi} + \mathbf{e}_3 \frac{\partial\psi}{\partial z} \\ \nabla \cdot \mathbf{A} &= \frac{1}{\rho} \frac{\partial}{\partial\rho} (\rho A_1) + \frac{1}{\rho} \frac{\partial A_2}{\partial\phi} + \frac{\partial A_3}{\partial z} \\ \nabla \times \mathbf{A} &= \mathbf{e}_1 \left( \frac{1}{\rho} \frac{\partial A_3}{\partial\phi} - \frac{\partial A_2}{\partial z} \right) + \mathbf{e}_2 \left( \frac{\partial A_1}{\partial z} - \frac{\partial A_3}{\partial\rho} \right) + \mathbf{e}_3 \frac{1}{\rho} \left( \frac{\partial}{\partial\rho} (\rho A_2) - \frac{\partial A_1}{\partial\phi} \right) \\ \nabla^2\psi &= \frac{1}{\rho} \frac{\partial}{\partial\rho} \left( \rho \frac{\partial\psi}{\partial\rho} \right) + \frac{1}{\rho^2} \frac{\partial^2\psi}{\partial\phi^2} + \frac{\partial^2\psi}{\partial z^2}\end{aligned}$$

Spherical  
( $r, \theta, \phi$ )

$$\begin{aligned}\nabla\psi &= \mathbf{e}_1 \frac{\partial\psi}{\partial r} + \mathbf{e}_2 \frac{1}{r} \frac{\partial\psi}{\partial\theta} + \mathbf{e}_3 \frac{1}{r \sin\theta} \frac{\partial\psi}{\partial\phi} \\ \nabla \cdot \mathbf{A} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_1) + \frac{1}{r \sin\theta} \frac{\partial}{\partial\theta} (\sin\theta A_2) + \frac{1}{r \sin\theta} \frac{\partial A_3}{\partial\phi} \\ \nabla \times \mathbf{A} &= \mathbf{e}_1 \frac{1}{r \sin\theta} \left[ \frac{\partial}{\partial\theta} (\sin\theta A_3) - \frac{\partial A_2}{\partial\phi} \right] \\ &\quad + \mathbf{e}_2 \left[ \frac{1}{r \sin\theta} \frac{\partial A_1}{\partial\phi} - \frac{1}{r} \frac{\partial}{\partial r} (r A_3) \right] + \mathbf{e}_3 \frac{1}{r} \left[ \frac{\partial}{\partial r} (r A_2) - \frac{\partial A_1}{\partial\theta} \right] \\ \nabla^2\psi &= \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial\psi}{\partial r} \right) + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial\theta} \left( \sin\theta \frac{\partial\psi}{\partial\theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2\psi}{\partial\phi^2} \\ &\quad \left[ \text{Note that } \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial\psi}{\partial r} \right) \equiv \frac{1}{r} \frac{\partial^2}{\partial r^2} (r\psi). \right]\end{aligned}$$